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Synergetic effect of the interface electric field and the plasmon electromagnetic field in Au-Ag alloy mediated Z-type heterostructure for photocatalytic hydrogen production and CO₂ reduction

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ABSTRACT

The challenge of synergistically optimizing different mechanisms limits the further improvement of plasmon-mediated photocatalytic activities. In this work, the CdS/Au-Ag/B-TiO₂ photocatalyst, combining a heterojunction structure with localized surface plasmon resonance (LSPR), is prepared by a high-temperature hydrogen reduction and hot solvent method. The novel synergistic interaction at the interface between the electric field and the plasmon electromagnetic field drives effective charge separation with improved photocatalytic efficiency. The Cd3Ti2 catalyst exhibits the highest photocatalytic activity, and the H₂ generation efficiency under the full solar spectrum is 15.97 mmol·h⁻¹·g⁻¹ with good stability as high as 81.3% in 24 h cycles. The CH₄ and CO precipitation performances are achieved at 14.2 μ mol·h⁻¹·g⁻¹ and 113.9 μ mol·h⁻¹·g⁻¹, respectively, in the CO₂RR process. The design of the synergistic mechanism of the dual electric field offers an important development in the field of plasma-mediated photoredox catalysis.

1. Introduction

Effective solar energy harvesting and conversion is a potential solution to the energy crisis [1]. Among various technologies, plasmonic photocatalysts have received significant attention, and research on localized surface plasmon resonance (LSPR) has achieved unprecedented development for enhancing light absorption in a wide spectral range. In plasma-mediated systems, the effects of the LSPR excitation from metal nanostructures supported on the semiconductor are divided into three types, including plasmon-induced direct electron transfer (DET) [2], local electromagnetic field enhancement (LEMF), and plasmon-induced resonance energy transfer (RET) [3,4]. Although some

mechanisms of plasmon-mediated photoredox catalysis have been recognized, it is still a challenge to develop the synergy effects between different mechanisms to achieve the organic integration of multiple functions.

Au and Ag appear in many studies as co-catalysts or plasma components. Ag or Au nanoparticles have been used for high-efficiency photocatalytic hydrogen production as a co-catalyst. Yu et al. rationally designed a core-shell $\rm TiO_2/Au@NiS_{1+x}$ structure with an $\rm H_2$ production rate of 9616 µmol·h $^{-1}$ ·g $^{-1}$ and an apparent quantum efficiency of 46% at 365 nm [5]. Wang et al. reported a novel Z-scheme system based on the Au decorated 5,10,15,20-tetrakis (4-trimethylammoniophenyl) porphyrin tetra (p-toluene sulfonate) functionalized iron-doped

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carbon nitride. The H₂ evolution activity of 3172.2 µmol·h⁻¹·g⁻¹ under UV-Vis irradiation was achieved by promoting the surface redox activities via efficient charge separation and transfer [6]. Meanwhile, amorphous AgSex was incorporated in situ onto metallic Ag as a novel and excellent H2-evolution active site to boost the interfacial H2-generation of Ag nanoparticles in a TiO2/Ag system. The as-prepared TiO₂/Ag@AgSe_x photocatalyst exhibited the H₂-production rate of 853 μ mol·h⁻¹·g⁻¹ [7]. During the LSPR excitation, some hot electrons in plasma may gain enough energy to inject into the conduction band of the adjacent semiconductor. Such a process is called plasmon-induced DET [8]. Franco et al. have experimentally unraveled the ultrafast photophysics of CdSe nanorods (NRs) and CdSe-Au HNPs, using a 2DES to describe the two-step relaxation of hot carriers [9]. In the work of Chen et al., the Au/XS2 core@shell nanostructures with well-controlled multilavers of the shell [10]. The unique Au/XS₂/Au antenna-reactor hybrids were synthesized using a hydrothermal method for the first time, which enhanced the transfer of hot electrons with improved photocatalytic hydrogen production performance up to 51.6 µmol·g⁻¹·h⁻¹. Li et al. realized that plasma hot electrons from oxygen vacancies could be used for the infrared-driven catalytic CO₂ reduction on Bi₂O_{3-x}, which enabled the efficient photocatalysis of CO₂ to CO (~100% selectivity, 16.15 μ mol·g⁻¹·h⁻¹) [11]. Dhiman et al. showed the tuning of the catalytic behavior of dendritic plasmonic colloidosomes (DPCs) in plasmonic hotspots, which successfully converted CO2 to methane (1.5 μ mol·g⁻¹·h⁻¹) at atmospheric pressure and room temperature, driven by sunlight [12]. Li's research team used H2O as an electron donor to achieve unassisted and highly selective gas phase CO2 reduction on a plasma-driven Au/p-GaN photocatalyst using photo-excited hot carriers. Its CO precipitation performance reached $\overline{650}~\mu\text{mol}\cdot\text{g}^{-1}\cdot\text{h}^{-1}$ [13]. Sistani et al. researched plasma-driven hot electron transfer at atomic-sharp metal-semiconductor nanojunctions, and the ability of momentum matching was demonstrated by controlling the energy distribution of the injected plasmon-driven hot electron [14].

Apart from the DET effect, the LEMF and RET effects are also important. LEMF can significantly boost the interband transition rate in the adjacent semiconductor if the plasmon resonance energy is larger than the semiconductor bandgap. On the other hand, the plasmonic oscillation energy could be transferred in the RET to the semiconductor photocatalysts through the EM field or non-radiative dipole-dipole coupling, which requires spectral overlap between the metal LSPR band and the light absorption range of the semiconductor [8,15,16]. The surface plasmon may improve the electron-hole pair generation and separation by these two mechanisms. Ma et al. prepared a core-shell (Au/AgAu)@CdS hybrid nanostructure with strong multi-interfacial plasmon coupling. The strong plasmon absorption and near-field enhancement can boost the generation and separation of electron-hole pairs in the multigap photocatalysts [17]. Yang et al. developed a dual-plasmonic-antenna strategy that efficiently generates energetic hot electrons with strong electromagnetic fields to boost photocatalytic hydrogen evolution (191.2 mmol \cdot g⁻¹·h⁻¹) [18]. Wang et al. used an asymmetric synthetic method to prepare Au-CdSe core-semi-shell nanorods for plasmon near-field-enhanced visible-light-driven hydrogen evolution (with a yield of 262.3 μ mol·g⁻¹·h⁻¹) [19]. Wang et al. have developed a facile surfactant-free nanofabrication approach to embed Au nanoparticles in N-doped TiO2 nanobowl arrays. The utilization of plasmonics in charge generation and separation by RET was well synergized with the Schottky junctions for enhanced photocatalytic water splitting, achieving the H₂ production rates of 637 μmol·g⁻¹·h⁻¹ illuminated with the full solar spectrum and 132 µmol·g⁻¹·h⁻¹ with visible light only [20]. Gao et al. revealed the polarization effect in surface plasmon-Induced photocatalysis on Au/TiO₂ nanoparticles [21]. Vahidzadeh et al. reported that the highest electric field intensities were found between the Ag-TiO2 and Ag-air interfaces. Under AM1.5 illumination, the highest rate of CH₄ production of 9.38 μ mol·g⁻¹·h⁻¹ was achieved on the AgCu-TNTA [22].

LSPR is a promising strategy for improving solar energy conversion

in plasmonic-enhanced photocatalysis. However, achieving efficient charge transfer with effective charge separation in metal/semiconductor systems is still challenging. In some cases, the near-field of isolated plasmonic nanoparticles cannot efficiently promote the charge dynamics in adjacent semiconductor photocatalysts. One possible way to solve this issue is the combination of two different types of metal nanoparticles in the same nanostructure to establish bimetallic nanoparticles (BNPs) [23,24]. Compared with isolated single metal nanoparticles (MNPs), BNPs exhibit a much stronger local plasmonic effect with a highly strengthened EM field due to their synergetic effect of its constituent elements at the metal-semiconductor interface. It has been reported that the electromagnetic field intensity of localized "hot spots" can reach 1000-fold of the far-field excitation light, electromagnetic field intensity based on the finite-difference time-domain (FDTD) simulations with discrete dipole approximation (DDA) methods [8,21,25]. Among many metals with LSPR effect, Ag offers excellent performance. However, its chemical stability limits its application. Au has better stability than Ag, but its LSPR strength is weak. Au-Ag alloys often show better enhanced catalytic activity than the isolated Au or Ag particles due to the stronger LSPR and higher stability under photoexcitation [26, 27]. Au-Ag alloy-mediated Z-type heterostructure can accelerate electron migration with improved electron-hole separation due to the synergetic effect of the electric field at the interface. The plasmonic electromagnetic field helps to reduce the random migration of low-energy photoexcited electrons and holes in the electronic medium and drives high-energy electrons to participate in the photocatalytic reactions, such as carbon dioxide reduction or hydrogen production.

In this work, we propose combining the LSPR with a heterogeneous electronic structure to achieve efficient charge transfer and separation. It is based on constructing the synergistic dual electric fields between the interfacial electric field and the plasmon electromagnetic field, which improves the vector charge transfer kinetics, resulting in a much improved photocatalytic carbon dioxide reduction performance with an optimal CH₄ production rate of 14.2 μ mol·g⁻¹·h⁻¹ and HER performance for an H₂ production rate of 15.97 mmol·g⁻¹·h⁻¹). First, we synthesized the TiO₂(B) with a standard hydrothermal method, which was reduced in an H₂ environment at high temperature to B-TiO₂ with a broadened light absorption wavelength range from visible to near-infrared. Au-Ag nanoparticles were synthesized separately in a standard method, which was impregnated onto the reduced B-TiO2 to form Au-Ag/B-TiO2 composite photocatalysts with different concentrations of Au-Ag. While, Since the plasmonic bands of Au-Ag alloys nanoparticles are mainly located in the visible wavelength range, the coupling between Au-Ag and B-TiO2 can effectively improve solar energy utilization through the RET mechanism. Hence, light absorption is no longer limited by the semiconductor bandgap energy. The finite-difference time domain (FDTD) methods represent a reasonable foundation for the LEMF mechanism. Since the three plasma energy transfer mechanisms are not mutually exclusive, it is expected that the RET mechanism and the LEMF mechanism will play major roles in the performance enhancement of the catalytic system [28,29]. In addition, when constructing the Au₃Ag₇--TiO₂ and Au₃Ag₇-CdS dielectric surrounding, the interfacial electric field can directly improve the vector charge transfer process, driving photoelectrons transfer from the semiconductors to the metallic NPs, which is verified by the DFT calculations. The band energy levels of CdS match well with those of TiO2. The existence of the dual-phase of TiO2 (anatase/TiO2(B)) also enables the CdS/Au-Ag/anatase/TiO2(B) system to form both II-type and Z-type composite heterojunctions. Effective charge separation will allow more electrons to participate in photocatalytic redox reactions [30]. The collaboration between the electric field at the interface and the plasma electromagnetic field offers the possibility of a fully optimized photocatalyst with an extended wavelength of light absorption and improved charge separation, which will become an important field in plasma-mediated photoredox catalysis.

2. Experimental section

2.1. Materials

Ethylene glycol ((CH₂OH)₂), titanium tetrachloride (TiCl₄), silver nitrate (AgNO₃), sodium borohydride (NaBH₄), trisodium citrate dihydrate (C₆H₅Na₃O₇), chloroauric acid (HAuCl₄·4 H₂O), cadmium nitrate (Cd(NO₃)₂·6 H₂O) and thiourea (CH₄N₂S) were purchased from Shanghai Aladdin Biochemical Technology Co., Ltd. Ethylenediamine (C₂H₈N₂) is supplied by Tianjin Shengao Chemical Reagent Co., Ltd. All chemicals are analytical grade, and used without further purification.

2.2. Synthesis of TiO₂(B) and B-TiO₂

First, based on our previous work [31], 60 mL ethylene glycol was added to an 80 mL PTFE-lined stainless steel hydrothermal reactor. Then, 1 mL TiCl₄ was added to the suspension until no HCl gas was formed at room temperature. After that, 1 mL of deionized water was added to the mixture. The sealed autoclave was heated in an oven at 150 °C for 4 h. Finally, the TiO₂(B) nanoflower product was collected by centrifuge. The impurities were removed by washing the solid with distilled water and absolute ethanol. The cleaned sample was vacuumed and dried for 24 h at 60 °C. The sample TiO₂(B) was calcined at 350 °C, 450 °C, and 550 °C in H₂/Ar (5%/95%) atmosphere to obtain mixed phase B-TiO₂ with different components (350-HB, 450-HB, and 550-HB). The prepared TiO₂(B) was added to a corundum crucible and heated in a muffle furnace at 550 °C with a heating rate of 5 °C/min for 2 h before the anatase sample was collected (550-AirB).

2.3. Synthesis of Au-Ag/B-TiO2

Au-Ag alloy NPs were deposited on the B-TiO $_2$ by an impregnation method. The detailed method for producing Au-Ag alloy is shown in the supplementary materials. For impregnation, 10 mL Au-Ag NP dispersion was added to a tube containing 50 mg B-TiO $_2$. The dispersion solution was thoroughly mixed using an ultrasonic disperser and then placed in a drying oven at 60 °C. The dried sample was collected and stored in the dark.

2.4. Synthesis of CdS/Au-Ag/B-TiO₂

Five CdS/Au-Ag/B-TiO₂ samples (Cd1Ti4, Cd2Ti3, Cd1Ti1, Cd3Ti2, and Cd4Ti1) were synthesized with different weight ratios of 1:4, 2:3, 1:1, 3:2, and 4:1 between CdS and Au-Ag/B-TiO₂. Cd(NO₃)₂·6 H₂O (0.966, 1.287, 1.159, 1.159, and 1.545 g) was mixed with the corresponding Au-Ag/B-TiO₂ (1, 0.5, 0.3, 0.2, and 0.1 g) precursor and CH₄N₂S (0.711, 0.949, 0.854, 0.854, and 1.138 g) in 30 mL ethylenediamine respectively. The solution was transferred to a 100 mL PTFE-lined autoclave and heated at 180 °C for 24 h. After the reaction, the sample was cooled to room temperature before being washed with ethanol and distilled water several times. The samples were vacuum-dried at 60 °C overnight. Pure CdS nanowires were synthesized following the same hydrothermal conditions without adding Au-Ag/B-TiO₂.

2.5. Characterization

Bruker D8 ADVANCE A25X X-ray diffractometer (AXS Co., Ltd.) was used to analyze the crystal structures of the samples. BET specific surface areas were measured with a QDS-MP-30 surface area analyzer. The microstructures were observed on a scanning electron microscope (SEM, Inspect F50, FEI) and a transmission electron microscope (TEM, JEM-2100 F. JEOL). The ultraviolet-visible (UV-Vis) diffuse reflectance spectra (UV-vis-DRS) were collected on a UV-Vis spectrophotometer (UV-2600, Shimadzu). X-ray photoelectron spectroscopy (XPS, ESCALA 250, Thermal Fisher) was used to analyze the elemental composition

and oxidation status. All binding energies were calibrated using the C 1 s peak at 284.8 eV. The photoluminescence (PL) spectrum was measured with a fluorescence spectrophotometer (LS55, Perkin Elmer) under the excitation wavelength of 315 nm. Electron paramagnetic resonance (EPR) spectra were measured at 77 k with an electron paramagnetic resonance spectrometer (JES-FA200, JEOL). Photoelectrochemical (PEC) measurements were performed on an electrochemical workstation (CHI 660E, Shanghai Chenhua Instrument Company). The photocatalytic reaction performances (HER and CO₂RR) were evaluated on a closed Pyrex gas circulation system (LabSolar 6 A, Beijing Perfectlight Co., Ltd.) at the atmospheric pressure with an irradiation cell at the top with the details in the supplementary materials.

2.6. FDTD simulations

The three-dimensional FDTD method was used to simulate the near-field distributions of the coupled nanoparticles. CdS/Au- Ag/TiO_2 materials were illuminated by a plane wave in the simulation with an excitation wavelength of 520 nm. Dielectric properties of different materials were obtained from the DFT calculations. The initial sizes and shapes of Au-Ag alloy, CdS, and TiO_2 in the simulation were estimated from the SEM and TEM images, as shown in Fig. 9. The Au-Ag alloy hollow nanospheres have a radius of 20 nm and a shell thickness of 10 nm. The CdS nanowire has a diameter of about 40 nm, simulated by a cylinder with a radius of 20 nm and a length of 70 nm. Due to the irregular morphology of TiO_2 , a bulk crystal model was used in the modeling. For easier comparison of electromagnetic strength, its modeling is consistent with CdS.

2.7. DFT calculations

First-principles calculations were carried out to investigate the properties of CdS/Au₃Ag₇/TiO₂ using the Vienna Ab-initio Simulation Package (VASP) [32]. The interaction between ions and valence electron densities is described using the projector augmented wave (PAW) potentials, and GGA-PBE is selected to treat the exchange correlation [33]. The cutoff energy of 450 eV was used for all calculations. The k-point mesh of $2\times2\times1$ was used. The relaxation of the electronic degrees of freedom is assumed to be converged when the total energy changes between the two electronic optimization steps are smaller than 1×10^{-5} eV, and the force below 0.05 eV/Å for ions was used as the criterion for relaxation convergence. The van der Waals effects were also considered with Grimme's DFT-D3 method [34,35].

3. Results and discussion

3.1. Structure characterization of the photocatalysts

The crystal structure, composition, and crystalline phases of the synthesized samples were characterized by X-ray diffraction (XRD). The XRD patterns of the pure TiO₂(B) and B-TiO₂ are shown in Fig. 1(a). The pure TiO2(B) diffraction peaks can be indexed to the monoclinic titanium dioxide (JCPDS no. 46-1238). After high-temperature hydrogen reduction, the 350-HB, 450-HB, and 550-HB samples show the three major peaks at 25.3°, 37.8°, and 48°, which are assigned to the (101), (004), and (200) crystal planes of the anatase phase (JCPDS no. 21-1272) [36]. The diffraction peak intensity of the anatase phase increases gradually with the increase in temperature, which indicates that crystallinity was improved as the reduction temperature was increased [37]. Therefore, the sample reduced at 550 °C (550-HB) was selected as the substrate for the subsequent deposition of Au-Ag and CdS. Interestingly, at the same temperature of 550 °C, the XRD pattern of 550-HB is different from that of the sample annealed in the air (550-AirB). The peaks of (103), (112), (105), and (211) were not observed in the 550-HB. Meanwhile, the peaks in 550-HB have larger widths than 550-AirB, demonstrating some amorphous structure was formed in the 550-HB

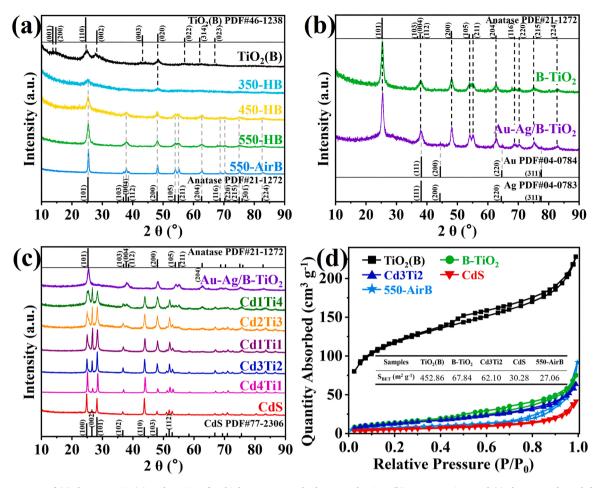


Fig. 1. XRD patterns of (a) the pure $TiO_2(B)$ and $B-TiO_2$ after high-temperature hydrogen reduction, (b) Au-Ag/B-TiO₂, and (c) the pure CdS and the composite samples. (d) N_2 adsorption-desorption isotherms for $TiO_2(B)$, B-TiO₂, CdS, Cd3Ti2, and 550-AirB.

[38]. To further determine the phase structure of 550-HB, we compared the XRD patterns of pure $TiO_2(B)$ and 550-HB using finer crystallographic information as a reference. What can be seen is that the (001), (200), (201), (310), (402), (003), (610), (712), (621), and (024) crystal planes belonging to $TiO_2(B)$ can all be observed from the 550-HB (Fig. S1). Hence, the 550-HB has a mixed anatase phase and $TiO_2(B)$, with anatase being the main component [36,39].

Fig. 1(b) shows the XRD pattern of the Au-Ag/B-TiO₂. The diffraction peaks of the B-TiO2 remained unchanged before and after adding the Au-Ag alloy, indicating that the added Au-Ag alloy does not affect the crystal structure of the B-TiO₂. The additional diffraction peaks at 38.1°, 44.3°, 64.4°, and 77.5° are indexed to the (111), (200), (220), (311) crystal planes of Au (JCPDS no. 04-0784) and Ag (JCPDS no. 04-0783) [40], respectively. Au and Ag have similar lattice constants. Therefore Au-Ag alloy has diffraction peaks similar to those of Au and Ag, consistent with previous reports [24]. In the Au-Ag/B-TiO₂ samples, Au and Ag diffraction peaks are very weak, owing to the good dispersion of lower content of Au-Ag. Nevertheless, the formation of Au-Ag alloy was identified in Fig. S3 [41]. The XRD spectra of the CdS/Au-Ag/B-TiO₂ composite samples with different CdS contents are shown in Fig. 1(c). The pure CdS shows the characteristic diffraction peaks at 24.8°, 26.5°, 28.2°, 43.7°, and 47.9°, attributed to the (100), (002), (101), (110), and (103) crystal planes, respectively, of the hexagonal wurtzite CdS (JCPDS no. 77-2306). Additional two weak peaks at 36.7° and 51.9° of the CdS crystals in the composite samples assigned to the (102) and (112) planes of the wurtzite CdS. Meanwhile, the characteristic peaks of B-TiO₂ were also observed distinctly in the CdS/Au-Ag/B-TiO₂ samples. With the increase of CdS loading, the intensities of the CdS peaks were progressively enhanced, while the relative intensities of the B-TiO2 XRD signals

were decreased accordingly. The results indicate the coexistence of B-TiO₂ and CdS in the heterostructured composite [42].

The specific surface areas of the samples were investigated by BET gas-sorption measurements. Fig. 1(d) shows the N2 adsorptiondesorption isotherms for TiO₂(B), B-TiO₂, CdS, and Cd3Ti2. The Cd3Ti2 was selected as the comparison sample since it offered the best photocatalytic and electrochemical performance. The nitrogen adsorption-desorption isotherms of TiO2(B), B-TiO2, and Cd3Ti2 show the typical type IV curves with the TiO2(B) and B-TiO2 showing the typical H3 hysteresis loops, meaning that the materials have slit holes formed by lamellae and particles, which was confirmed by the subsequent morphological studies (Fig. S3) [43]. The measured specific BET surface areas of the TiO2(B), B-TiO2, Cd3Ti2, and CdS were 452.86, 56.81, 62.1, and 23.63 m $^2 \cdot g^{-1}$, respectively. The lower surface area of the $B\text{-TiO}_2$ compared to the $TiO_2(B)$ was attributed to the transformation of TiO2(B) nanosheets into TiO2 nanoparticles during the high-temperature hydrogen reduction. The CdS shows a type II physisorption isotherm curve. Because of this, the void structure of the sample Cd3Ti2, composed of B-TiO2 and CdS, has the structural features of both components with less distinctive of the H3-type hysteresis loop. The capillary condensation brought by the slit pores is beneficial to the adsorption of reactants. A heterostructured photocatalyst with a large specific surface area offers a high density of active sites and facilitates photo-generated carrier transport for redox reactions, making it favorable for enhancing photocatalytic performance [42,44].

The morphology and microstructure of Au-Ag alloy, $TiO_2(B)$, B- TiO_2 , CdS, and CdS/Au-Ag/B- TiO_2 were studied by SEM and TEM. The SEM images of the Au-Ag alloy are shown in Fig. S3 (a-c), with an average particle size of \sim 40 nm. A film structure on the particle surface may be

due to the residual citrate ligand. TEM (Fig. 3(a) and (b), Fig. S2 (c) and (d)) images identified the hollow nanoparticulate structure of the Au-Ag alloy. Fig. S3 (d-f) shows the SEM images of the TiO₂(B) with a spherical nanoflower morphology at a dimension of 200-500 nm formed by stacking nanosheets, consistent with previous reports. Hence, it is expected to have a high specific surface area, as confirmed by the BET results in Fig. 1(d). After high-temperature calcination in H₂, the nanosheet structure of TiO₂(B) changed to fine particles with irregular spherical B-TiO₂ structure (200-500 nm) (Fig. S3 (g-i)), formed by the accumulation of small particles in the range of 20-50 nm. Meanwhile, Au-Ag alloy was successfully deposited on the B-TiO2, as shown in Fig. 2 (a-d), demonstrated by some small particles attached to the surface of B-TiO₂, forming the Au-Ag/B-TiO₂ heterojunction [45]. Pure CdS shows a regular nanowire morphology with a relatively smooth surface. The average length is ~5 µm with a diameter of 40 nm, measured from the images in Fig. 2(e-g). Fig. 2(h-j) show the SEM images of the best-performing composite sample, Cd3Ti2. There is no obvious damage to the Au-Ag/B-TiO2 structure in the solvothermal synthesis of CdS nanowires. Fig. S4 (a-f) show the SEM images of CdS/Au-Ag/B-TiO₂ samples with a different CdS content, which show the significant structural changes of the CdS nanowires due to the addition of Au-Ag/B-TiO2. It should be noted that with the increase of added Au-Ag/B-TiO2, the CdS nanowires become shorter. Hence, Au-Ag/B--TiO₂ could inhibit the growth of CdS nanowires along the axial direction

The surface structure and morphology of Au-Ag alloy and Cd3Ti2 were further examined with TEM. Fig. 3(a) and (b) show that the Au-Ag alloy has a uniform hollow nanoparticulate structure with a size of 40–50 nm and a shell thickness of 20 nm. The lattice fringes of Au-Ag alloy give the lattice spacings of 0.235 and 0.204 nm observed in a high-resolution transmission electron microscope (HRTEM) could be

indexed to the (111) and (200) crystal planes of the Au-Ag alloy (Fig. 3 (c)). Besides, energy-dispersive X-ray spectroscopy (EDS) mapping analysis shows a homogeneous distributed Au and Ag in the Au-Ag alloy, which further confirms that the Au-Ag solid solution was successfully prepared at the Au: Ag atomic ratio of 3:7.

The TEM images in Fig. S5 reveal that the Au-Ag alloy nanoparticles were evenly dispersed on the TiO₂ surface. To demonstrate that Au-Ag alloy nanoparticles have good interaction with the TiO₂ surface, the structures of Au-Ag/TiO₂(B) (Fig. S5 (a-c)) and Au-Ag/B-TiO₂ (Fig. S5 (d-f)) were investigated [47]. Au-Ag alloy nanoparticles were deposited via a dipping method. They are all contacting interfaces, and there exists an interaction between Au-Ag alloy and TiO₂. Here, we speculate that a certain amount of citrate ions were left on the surface of the Au-Ag alloy during the synthesis, which helps to interact with the oxygen vacancies (Ovs) or surface OH groups on the TiO₂(B) and B-TiO₂ surfaces. The direct coupling of Au-Ag NPs and B-TiO₂ is conducive to transporting electrons between two materials [45]. Further observation of the lattice fringes of the Au-Ag/B-TiO₂ shows the lattice spacings of 0.240, 0.352, and 0.358 nm, corresponding to the (001), (101) planes of anatase and the (110) plane of TiO₂(B), respectively (Fig. 3(d)).

The TEM images further confirm the coupling of Au-Ag NPs, B-TiO $_2$, and CdS in CdS/Au-Ag/B-TiO $_2$. Fig. 3(i-j) show the CdS crystals in the sample Cd3Ti2 having a nanowire morphology. Au-Ag NPs, as the intermediate medium, are sandwiched between B-TiO $_2$ and CdS nanowires. The well-defined boundaries between the Au-Ag/B-TiO $_2$ and CdS nanowires indicate the successful formation of the heterojunctions in the composite (Fig. 3(k)). A magnified area of the CdS/Au-Ag/B-TiO $_2$ heterojunction is shown in Fig. 3(l). The catalyst shows a clear lattice fringe with the corresponding fringe spacing of 0.336 nm, ascribed to the (002) facet of the CdS nanowires, consistent with the XRD results [48]. Element mapping results from EDS (Fig. 3(m-o)) verified the close

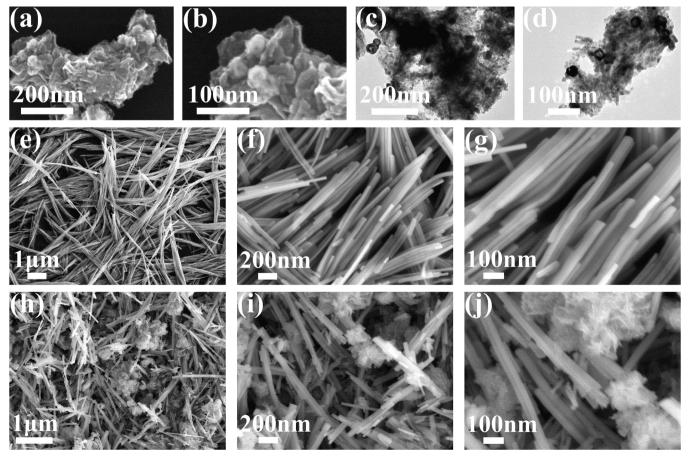


Fig. 2. The morphology images of (a-d) Au-Ag/B-TiO₂. SEM images of (e-g) the CdS nanowires, and (h-j) the sample CdS/Au-Ag/B-TiO₂.

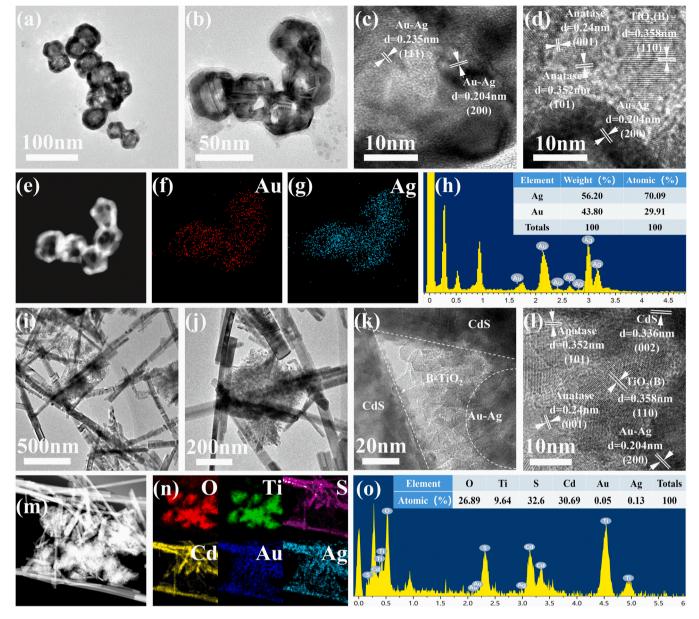


Fig. 3. (a), (b) TEM images of the Au-Ag alloy. HRTEM images of (c) the Au-Ag alloy, and (d) the Au-Ag/B-TiO₂. (e-h) Images of HAADF-STEM and EDS mapping for the Au-Ag alloy. (i) (j) TEM images of the Cd3Ti2. (k) (l) HRTEM images of the Cd3Ti2. (m-o) Merged images of HAADF-STEM and EDS mapping for the Cd3Ti2.

contact between the Au-Ag NPs and the B-TiO₂ or the CdS nanowires. Meanwhile, the O, Ti, S, Cd, Au, and Ag were uniformly distributed in the Cd3Ti2, confirming the formation of the CdS/Au-Ag/B-TiO₂ heterojunction [38]. Therefore, the Au-Ag NPs played a significant role in determining the mobility of the charge carriers, which is responsible for the improved photocatalytic activity of CdS/Au-Ag/B-TiO₂ [41].

3.2. Energy band structure and mechanism

Optical properties are important in influencing photocatalytic performance. The optical absorption spectra were measured by UV-Vis DRS to determine the band gaps, as shown in Fig. 4(a). The absorption edge of ${\rm TiO_2}(B)$ was around 375 nm, while that of anatase was around 420 nm. The corresponding band gaps were 3.6 eV and 3.2 eV, respectively (Fig. 4(b)), which agrees with previously reported values. Hence, ${\rm TiO_2}$ samples only exhibit a UV light response. A redshift of the absorption edge of B-TiO₂ is revealed after the high-temperature hydrogen reduction of ${\rm TiO_2}$. The color of B-TiO₂ changed from white to black

(Fig. 4(d)), and visible light responsiveness was significantly enhanced due to the presence of oxygen vacancies and ${\rm Ti}^{3+}$. The small peaks at 540 and 675 nm of Au-Ag/B-TiO₂ in the inset may be caused by the LSPR effect of the Au-Ag NPs, assigned to the quadrupole and dipole plasmon signals [18,49]. These characteristic peaks are not particularly obvious due to the low content of the Au-Ag NPs. Fig. 4(a) shows a broadened and slightly enhanced light absorption edge of Au-Ag/B-TiO₂ compared with B-TiO₂, indicating that the Au-Ag NPs form additional states within the band gap [50], effectively reduce the band gap, and improve the absorption of visible light.

The additional states can also promote charge transfer. The electrons at the top of the valence band can transfer to the Au-Ag state within the band gap [48]. Since the plasmonic bands of the Au-Ag NPs are mainly located in the vis-spectral range overlapped with the B-TiO₂, the RET mechanism can be valid [51,52], resulting in enhanced visible light absorption, charge carrier generation, and effective charge separation [45]. Therefore, the combination of Au-Ag/B-TiO₂ can significantly enhance the utilization of visible light and generate more charge

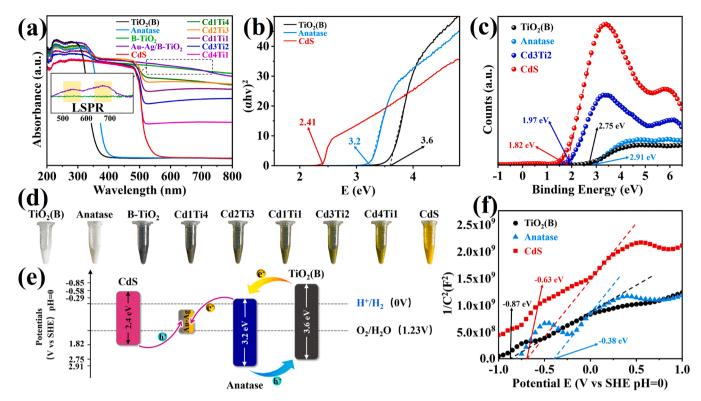


Fig. 4. (a) UV-Vis absorption spectra of TiO₂(B), anatase, B-TiO₂, Au-Ag/B-TiO₂, CdS and composite catalysts CdS/Au-Ag/B-TiO₂, and (b) the absorption spectra from Kubelka-Munk transformed functions. (c) XPS valence band spectra of TiO₂(B), anatase, Cd3Ti2 and CdS. (d) Color of the samples. (e) The electronic band structure of CdS/Au-Ag/B-TiO₂. (f) Mott-Schottky plots of TiO₂(B), anatase and CdS.

carriers.

Fig. 4(a) shows the absorption edge of pure CdS nanowires at approximately 530 nm with a corresponding bandgap of 2.41 eV. By coupling CdS with Au-Ag/B-TiO₂, the light absorption edges of composite catalysts, Cd1Ti4, Cd2Ti3, Cd1Ti1, Cd3Ti2, and Cd4Ti1, show regular changes. It is worth noting that the absorption edges of the composite catalysts are larger than those of B-TiO₂ and CdS. The results showed that combining B-TiO₂ and CdS could compensate for each other to improve their visible light absorption [39].

To investigate the valence band edge potentials of the catalysts and understand the influence of the band structure on the photocatalytic activity, the valence band edges of TiO2(B), Anatase, and CdS were determined by XPS valence band spectroscopy. Fig. 4(c) shows the valence band maxima (VBM) of TiO₂(B), anatase, and CdS catalysts that occurred at 2.75 eV, 2.91 eV, and 1.82 eV, respectively. Hence, the corresponding conduction band minima of TiO2(B), anatase, and CdS catalysts are determined to be -0.85, -0.29, and -0.58 V, respectively [48]. The electronic band structure alignments are presented in Fig. 4(e). It shows the influence of introducing Au-Ag NPs on the energy band structure of TiO2(B), anatase, and CdS catalysts. The CdS/Au-Ag/B-TiO2 composite catalyst exhibits a tandem of type II and type Z heterojunction. The Au-Ag NPs act as an electron-transporting medium. UV and partially visible light can be absorbed and excite the B-TiO₂, containing both TiO₂(B) and anatase, to generate the electron-hole pairs. The separation of these excited charges is improved by the type II heterojunction formed between TiO₂(B) and Anatase. Meanwhile, CdS can absorb visible light to produce excited electrons and holes. The holes in the CdS can recombine with the excited electrons transferred from the anatase through Au-Ag NPs while retaining its high-energy electrons on the conduction band of the CdS for the photocatalytic reaction. In addition, the visible light irradiation activates the surface plasmon of Au-Ag NPs. The RET process can effectively promote the generation and separation of the electron-hole pairs, which is favorable for the HER and CO₂RR reactions under sunlight [20].

Mott-Schottky plots were also used to reveal the energy band structure by identifying the flat band potentials of $TiO_2(B)$, anatase and CdS catalysts to verify the possible photocatalytic mechanism [53]. The results are shown in Fig. 4(f). All samples were identified as n-type semiconductors due to the positive slopes in the Mott-Schottky plots. The CB of an n-type semiconductor is usually about 0.1 V below the flat band potential [43]. Consequently, the E_{fb} (vs. SHE) of $TiO_2(B)$, Anatase, and CdS were determined to be - 0.87, - 0.38, and - 0.63 V_{SHE} , which is close to the results from the XPS-VB and UV-Vis measurements [54, 55]. Therefore, the synergistic effect of the heterojunction electron transfer mechanism and the RET mechanism for improved visible light absorption can improve the efficiency of sunlight utilization to a certain extent, and provide reliable performance for the photocatalytic hydrogen evolution reaction and carbon dioxide reduction reaction.

3.3. Identification of chemical bond

To explore the surface composition and chemical states of elements and to confirm the charge transfer mechanism in the CdS/Au-Ag/B-TiO₂ heterojunction, XPS spectra of TiO₂(B), Anatase, Au-Ag/B-TiO₂, Cd3Ti2, and CdS were analyzed. The survey XPS spectral in Fig. 5(a) verifies the existence of Ti, O, Cd, S, Au, Ag, and C elements, in accord with the result of EDS. The atomic percentages of the elements are listed in Table S1. All samples are pure and free from other impurities. Fig. 5(f) displays the C1 s peaks at 284.8 eV, attributed to surface contaminated carbon, which was used to calibrate the binding energies of all elements. Fig. 5(b) shows the Au 4 f spectrum from Au-Ag/B-TiO₂ with two main peaks at 83.98 and 87.69 eV, assigned to the surface Au species [56]. The peaks at 367.95 and 374.02 eV in Fig. 5(c) are assigned to the Ag [57]. After loading CdS, the relative contents of Au and Ag in the composite sample of Cd3Ti2 decreased (Table S1), as shown in Fig. 5(b) and (c). Spectra from TiO₂(B), Anatase, Au-Ag/B-TiO₂, and Cd3Ti2 show symmetrical Ti 2p doublets from Ti⁴⁺ ions (Fig. 5(d)). Two additional peaks from the $TiO_2(B)$ at 463.42 and 457.98 eV are ascribed to the Ti^{3+}

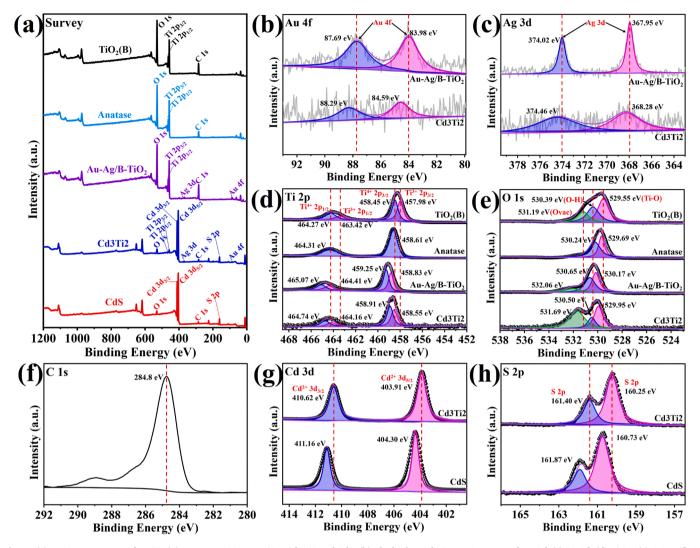


Fig. 5. (a) XPS survey spectra for TiO₂(B), anatase, Au-Ag/B-TiO₂, Cd3Ti2 and CdS. (b) The high resolution XPS spectra of Au 4 f, (c) Ag 3d, (d) Ti 2p, (e) O 1 s, (f) C 1 s, (g) Cd 3d, (h) S 2p.

 $2p_{1/2}$ and Ti^{3+} $2p_{3/2}$, respectively. The O 1 s spectrum from $TiO_2(B)$ is shown in Fig. 5(e), with three peaks at 529.55, 530.39, and 531.19 eV, which are assigned to the lattice oxygen (Ti-O), the adsorbed oxygen (O-H) [46], the abundant oxygen vacancies (Ovac) [45]. For the anatase obtained from calcining $TiO_2(B)$ at 550 °C in air, the Ovac peak disappeared in Fig. 5(d) and (e). Meanwhile, the hydrogen reduction at high temperatures inhibited the formation of defects-free anatase while retaining the density of oxygen defects.

For the Au-Ag/B-TiO₂ composite, the two Ti⁴⁺ 2p peaks are shifted to 465.07 and 459.25 eV, respectively. According to the previous report, the electron concentration can increase when the bond is formed between the two-phase interfaces [45]. Both the Ti 2p and the O 1s binding energies are shifted to higher binding energies compared with TiO₂(B) (Fig. 5(d) (e)). This indicates a reduced electron density in TiO₂(B), suggesting that the photo-generated electrons transfer from TiO₂(B) to anatase, and are further extracted by Au Ag alloy. At the same time, the Cd3Ti2 heterojunction composite displayed negative shiftings in the binding energies of ${\rm Ti}^{4+}$ 2p (-0.34 eV) and O 1 s (-0.22 eV), demonstrating that the multielement Cd3Ti2 heterojunction has the strong interaction, which is conducive the transfer of photoinduced carriers [39,48,58]. The Cd 3d binding energies from the composite Cd3Ti2 at 410.62 and 403.91 eV are assigned to the Cd $3d_{3/2}$ and $3d_{5/2}$, respectively (Fig. 5(g)). These peaks are redshifted to the lower binding energies than Cd 3d in CdS (411.16 and 404.30 eV). Meanwhile, the binding energy of S 2p shows similar shifts (Fig. 5(h)). Such shifts might be due to the difference in the work functions at the interface, causing charge transfer at the heterojunction interface, with the CdS as the electron acceptor [44].

Meanwhile, the binding energies of Au and Ag in their high-resolution XPS spectra in the Cd3Ti2 heterojunction shift towards the higher binding energies in comparison with that of Au-Ag/B-TiO₂, suggesting that the Au-Ag NPs are the electron donor. Due to the high CB position of the CdS (Fig. 4(e)), it is not conducive to the flow of the electrons from Au-Ag NPs to CdS, which indicates that the heterojunction system follows a Z mechanism. The result substantiates that the electrons migrate from the B-TiO₂ component to the Au-Ag alloy under the light irradiation, the low-energy photo-generated holes from CdS combine with low-energy photo-generated electrons gathered on the Au-Ag alloy, while the high-energy photo-generated electrons on CdS participate in the reaction, which is consistent with the proposed charge transfer mechanism [59].

3.4. Photoelectrochemical measurements

The photocatalytic properties of CdS/Au-Ag/B-TiO₂ composites are further explored by the transient photocurrent responses and electrochemical impedance spectroscopy (EIS). As shown in Fig. 6(a), The photocurrent response of pure TiO₂(B) is very weak, which is caused by

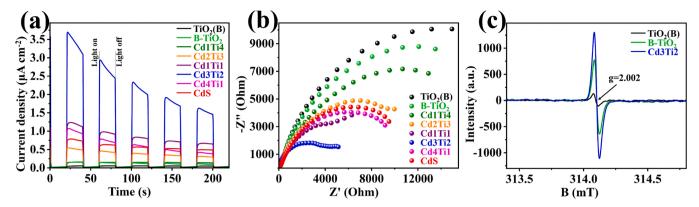


Fig. 6. (a) Transient photocurrent responses for TiO₂(B), B-TiO₂, CdS and composite catalysts CdS/Au-Ag/B-TiO₂, and (b) The electrochemical impedance spectra (EIS). (c) Electron paramagnetic resonance (EPR) of TiO₂(B), B-TiO₂ and Cd3Ti2.

the poor response to visible light and fast charge recombination. The photocurrent intensity of the CdS/Au-Ag/B-TiO $_2$ composites has been greatly improved. In particular, the average value of photocurrent intensity of 2.13 μ A cm $^{-2}$ on the Cd3Ti2, which is about 42 times that on the pristine TiO $_2$ (B) and 3.8 times that on CdS. Hence, the heterojunction structures of Cd3Ti2 can effectively separate photo-generated electron and hole pairs to improve photocatalytic performance [46]. In addition, the near-field electromagnetic coupling of the Au Ag plasma also plays an important role, which showed that the rate of electron-hole formation in semiconductors TiO $_2$ and CdS was enhanced in the regions near the Au Ag NPs with excited plasmonic effects [25].

Meanwhile, the electrochemical impedance spectroscopy (EIS) measurements are shown in Fig. 6(b) to illustrate the electron transfer process. The semicircle radius of Cd3Ti2 is smaller than those of other samples, implying less interface resistance for carrier transfer due to the enhanced electron mobility, which agrees with the photocurrent results [43,45]. All this data together identifies the fast electron transferring in Cd3Ti2, implying a promoted charge separation owing to the formation of a heterojunction [39,60].

To confirm the improved charge separation, the PL spectra of $TiO_2(B)$, CdS, and Cd3Ti2 in the range of 360–500 nm are collected and shown in Fig. S6. Since the PL emission results from charge recombination, a lower PL intensity generally indicates a higher charge separation efficiency [60]. The PL emission peak intensity of $TiO_2(B)$ is the highest in the testing samples, implying the highest electron-holes recombination rate. On the contrary, Cd3Ti2 has the lowest PL intensity, evidencing effective inhibition of the charge recombination, which is consistent with the results of the photocurrent responses and EIS measurements [43]. Therefore, the multiple heterojunction structures and the RET mechanism can efficiently promote the separation and transportation of electron-hole pairs and inhibit their recombination.

Electron paramagnetic resonance (EPR) has been widely used to reveal the existence of Ovac directly. As shown in Fig. 6(c), a significant EPR signal with a g value of 2.002 is observed in the samples of $TiO_2(B)$, B- TiO_2 , and Cd3Ti2, which indicates the existence of Ovac accompanied with the Ti^{3+} in these samples, which is in good agreement with the XPS results [38]. Further comparison in the intensities of the EPR signals revealed that the B- TiO_2 and Cd3Ti2 have higher Ovac concentrations than the $TiO_2(B)$, proving that high temperature hydrogen reduction can significantly promote the production of Ti^{3+} and Ovac. The existence of Ti^{3+} and Ovac is significant for improving photocatalytic performance.

3.5. Photocatalysis performance

Photocatalytic performance is evaluated to reveal the advantages of the catalytic system. Here, two of the most common photocatalytic reactions, the photocatalytic HER reaction and the photocatalytic CO_2RR reaction, were carried out. The photocatalytic hydrogen evolution

properties of CdS/Au-Ag/B-TiO2 composites were assessed under the full solar spectrum (AM=1.5). The photocatalytic hydrogen evolution performances of pristine TiO₂(B) and CdS were also examined under identical reaction conditions. The photocatalytic hydrogen evolution amount versus irradiation time over the CdS/Au-Ag/B-TiO₂, TiO₂(B), and CdS are plotted in Fig. 7(a). The amount of the produced H₂ was linearly increased with the reaction time. Accordingly, the reactivities of CdS/Au-Ag/B-TiO2 with different contents of CdS are better than TiO₂(B) and CdS. Results confirmed the superior activity of Cd3Ti2, achieving an H_2 production rate of 15.97 mmol·h⁻¹·g⁻¹, which is 24.2 times higher than that of the $TiO_2(B)$ (0.66 mmol·h⁻¹·g⁻¹), and 17.9 times higher than that of CdS (0.89 mmol·h $^{-1}$ ·g $^{-1}$), respectively (Fig. 7 (b)). The obtained apparent quantum efficiency (AQE) of Cd3Ti2 at 365 nm and 420 nm are 10.6% and 1.7%, respectively. The AQE of CdS, B-TiO₂, and Cd3Ti2 under different wavelengths of monochromatic light irradiation are summarized in Table S2.

The excellent HER performance of Cd3Ti2 could be attributed to the following aspects. Firstly, high-temperature hydrogen reduction improves the visible light absorption of TiO₂(B); secondly, the synergy of the type II heterojunction and the Z-type heterojunction greatly inhibits the recombination of the electron-hole pair; besides, the Au-Ag NPs between the TiO2 and CdS form a strong localized surface plasmon resonance (LSPR). Among them, the LEMF mechanism further enhances the charge transfer of the Z-type heterostructures. The RET mechanism also enhances the charge carrier generation and separation. Comparing related reports on the plasma mediated photocatalysis of hydrogen production (Table S3), the CdS/Au-Ag/B-TiO2 produced in this work offers the best HER performance. To further demonstrate the superiority of Au-Ag alloy, we replaced the Au-Ag alloy in the CdS/Au-Ag/B-TiO₂ with Ag nanoparticles. The photocatalytic performance of CdS/Au-Ag/ B-TiO₂ and CdS/Ag/B-TiO₂ are compared in Fig. S7. The Au-Ag alloy in the CdS/Au-Ag/B-TiO₂ system was more beneficial to the photocatalytic performance than the pure Ag nanoparticles. Cycling photocatalytic hydrogen evolution experiments were carried out to study the stability of the Cd3Ti2 HER performance [39]. As shown in Fig. 7(c), the H₂ production rate of Cd3Ti2 maintains fairly constant during 12 cycles of reaction, decreasing slightly to 81.3%, indicating good stability of Cd3Ti2 [43]. The Cd3Ti2 is subjected to morphological observation after long-cycle experiments, as shown in Fig. S8, and no significant changes in its structure occurred. We speculate that the photocorrosion of CdS is the main cause of the performance degradation in the Cd3Ti2.

The performances of the various photocatalysts for CO_2 reduction were evaluated under the full solar spectrum (AM=1.5). As shown in Fig. 8(a) and (b), the main products of the reaction are CH₄ and CO [61–63]. Both $TiO_2(B)$ and CdS showed low photocatalytic activities toward CO_2 reduction. Interestingly, the heterojunction photocatalyst of CdS/Au-Ag/B- TiO_2 showed a significantly improved activity for CO_2 reduction [46]. The photocatalytic activity can be optimized by

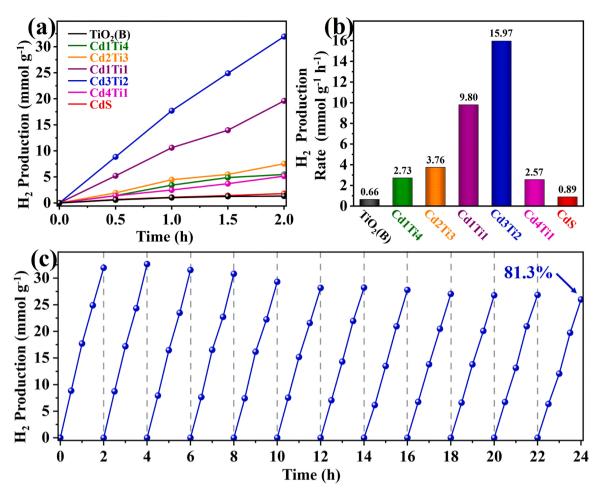


Fig. 7. Photocatalytic hydrogen evolution reaction (HER). (a) H₂ production under simulated solar irradiation of TiO₂(B), CdS and composite catalysts CdS/Au-Ag/B-TiO₂, and (b) Corresponding hydrogen production rates. (c) Stability test of Cd3Ti2 (the reaction system was evacuated for 20 min every 2 h).

adjusting the relative proportions of B-TiO₂ and CdS. As shown in Fig. 8 (c), the photocatalytic activities of the CdS/Au-Ag/B-TiO₂ show slight changes as the mass ratio of CdS to B-TiO₂ increases from 1:4-1:1. However, when the mass ratio of CdS to B-TiO₂ exceeded 1:1, the CO₂ reduction rates begin to improve significantly [58]. Among all composite photocatalysts, samples Cd3Ti2 and Cd4Ti1 exhibited the optimal CH₄ and CO production rates, respectively. The Cd3Ti2 shows the highest CH_4 yield rate of 14.2 μ mol·h⁻¹·g⁻¹, and is almost 15.8 times the yield from CdS (CH₄ production activity was not detected in TiO₂(B)). The Cd4Ti1 shows the highest CO yield of 327.2 μmol·h⁻¹·g⁻¹, almost 5.4 and 36 times the yields from TiO₂(B) and CdS, respectively [64]. The reactions producing CH₄ and CO involve 8 and 2 electrons, respectively. For the Cd3Ti2, the CH₄ content accounts for 11.1% (Fig. 8(d)), and CO is the main product of the CO2RR reaction. The literature results of plasma mediated photocatalytic CO2 reduction are compared in Table S4. According to the above analysis, the combination of Type II and Z-Type composite heterojunctions, RET and LEMF mechanisms has improved the directional charge transfer kinetics, resulting in the high performance of photocatalytic carbon dioxide reduction and HER.

3.6. Finite-difference time-domain (FDTD) simulations

The internal mechanism of the Au-Ag NPs LSPR effects on the photocatalytic performance of CdS/Au-Ag/TiO $_2$ is investigated using the FDTD simulation to study the LEMF mechanism, as shown in Fig. 9. A hot spot representing the focused electromagnetic field is identified in the junction under the excitation of a 520 nm light polarized along the X-axis. The electromagnetic intensity of the local hot spot is much higher

than that of the far-field excitation light, where the strength of the electromagnetic field is at least one order of magnitudes higher than that of the far-field excitation light [65]. The carrier generation and separation rates in the hot spot region will significantly increase. The results show that the Au-Ag NPs in the CdS/Au-Ag/TiO₂ effectively increase the electromagnetic field intensity, especially at the interfaces around the Au-Ag NPs interfaced with CdS or TiO2. Quantitatively, the electromagnetic field enhancements from the Au-Ag NPs, CdS/Au-Ag, and Au-Ag/TiO2 are 1.3, 9, and 8.6, respectively. This means that the electromagnetic fields for CdS/Au-Ag or Au-Ag/TiO2 are about 1 order of magnitude higher than that of Au-Ag NPs, indicating the importance of the heterojunction structure formed between Au-Ag NPs and the coupled semiconductors of CdS and TiO2. In conventional semiconductor-metal-semiconductor Z-type heterojunctions, low-energy electrons and low-energy holes recombine at the metallic components, thereby retaining high-energy electrons and holes for redox reactions. It is speculated that the strong electromagnetic field of the Au-Ag NPs will accelerate the transfer and combination of the low-energy electrons and holes in the electromagnetic field while suppressing the recombination of the high-energy electrons and holes, as confirmed by the EPR results. In summary, the intense plasmonic EM field at the metal nanostructures can significantly increase the interband transition rate in the adjacent semiconductors and facilitate the separation of photo-generated electron-hole pairs [18,22].

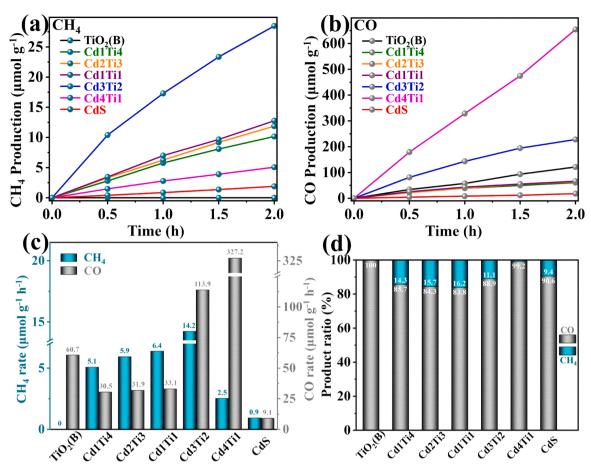


Fig. 8. Photocatalytic Carbon dioxide reduction reaction (CO₂RR). (a) CH₄ production under simulated solar irradiation of TiO₂(B), CdS and composite catalysts CdS/Au-Ag/B-TiO₂, and (b) CO production, and (c) Corresponding production rates of CH₄ and CO, and (d) The distribution of carbon dioxide reduction products.

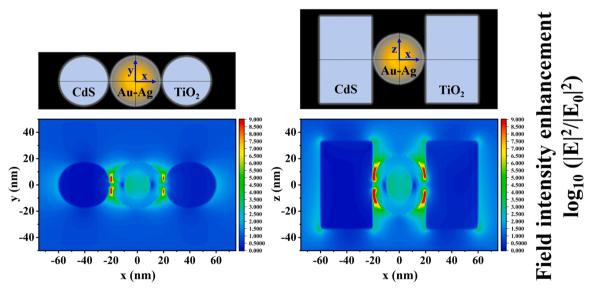


Fig. 9. Electric field profile for CdS/Au-Ag/ TiO_2 in the substrate (xy) plane and (xz) plane. The incident plane wave is polarized in the x-direction and propagates along the z-direction.

3.7. Density functional theory (DFT) calculations

3.7.1. Structure Information

Combining the elemental analysis (XPS and EDS) with our previous work [26,27], the element ratio of Au-Ag NPs can be obtained as 3:7. To

calculate the electronic structures of the CdS/Au₃Ag₇/TiO₂ photocatalyst, we first built a three-layers Ag (111) with 3×3 supercell, and contain 27 Ag atoms. To simulate the Au₃Ag₇, eight Ag is replaced by Au, distributed uniformly. The lattice parameter is 8.65, 8.65, and 29.71 Å. CdS (002) is built by six layers with 2×2 supercell, including a

vacuum layer of 15 Å with the lattice parameter of 8.41, 8.41, and 22.68 Å. Hence, the lattice mismatch ratio is less than 5%. A similar method was used for constructing the Au₃Ag₇-TiO₂ composite, which contains 48 Ti, 80 O, 20 Au, and 52 Ag, as shown in Fig. 10.

3.7.2. The electronic structure analysis

DFT calculations provide further insight into the underlying mechanisms. The work function (Φ) is a crucial parameter for electron transfer within heterojunctions, defined by the energy difference between the vacuum and Fermi levels. The work functions of Au_3Ag_7 -CdS and Au_3Ag_7 -TiO $_2$ were calculated to establish the possible migration within the CdS/ Au_3Ag_7 /TiO $_2$ material. As shown in Fig. 11 (a) and (d), the calculated work functions of the Au_3Ag_7 -TiO $_2$ and Au_3Ag_7 -CdS are 4.64 and 6.35 eV, respectively. Hence, the Au_3Ag_7 -CdS has a lower Fermi level than the Au_3Ag_7 -TiO $_2$.

Consequently, the electrons will transfer from the Au_3Ag_7 - TiO_2 to Au_3Ag_7 -CdS across their interface to find a balance of Fermi levels. Thus, a built-in electric field was formed at the interface between Au_3Ag_7 -CdS and Au_3Ag_7 -TiO $_2$, pointing from the Au_3Ag_7 -CdS to Au_3Ag_7 -TiO $_2$, which drives the migration of the photo-generated electrons and holes. Therefore, it is expected that, with the CdS/ Au_3Ag_7 /TiO $_2$ heterojunction excited under sunlight, the low-energy photo-generated charges will be transferred from the TiO $_2$ to the Au-Ag NPs. Conversely, the low-energy photo-generated holes will be transferred from the CdS to Ag-Au NPs. After their annihilation, the energetic electrons will be left on the CB of the CdS, while the excited holes will be left on the VB of the TiO $_2$. As

such, the photoexcited charges are well separated [4,66].

DFT calculations were also performed to identify further the electronic structure of the CdS/Au₃Ag₇/TiO₂. As shown in Fig. 11 (b) and (e), the electron density of the conduction band (CB) of Au₃Ag₇-TiO₂ decreases, whereas the electron density of CB of Au₃Ag₇-CdS increases (PDOS). This is revealed from the density of state (DOS) and partial density of state (Fig. S9) in Fig. 11 (b) and (e). Fig. S10 shows the electron depletion and accumulation in the Au₃Ag₇-TiO₂ and Au₃Ag₇-CdS. This proves the existence of strong interactions across the interfaces, eliminating some dangling bonds and generating interfacial electronic states. The electrostatic potential shows an abrupt change along the interface directly with a contact potential difference. In the Au₃Ag₇-CdS interface, the blue bubbles represent the dissipation of electrons, and the yellow bubbles represent the aggregation of electrons. The interfacial potential difference drives the transfer of photogenerated holes on CdS to the alloy region. Similarly, the opposite process exists at the Au₃Ag₇-TiO₂ interface. This can further enhance the Z-type heterojunction mechanism and improve photocatalytic activity. By forming the CdS/Au-Ag/B-TiO2 composite, a strong electrical field is expected to be formed by moving electron density from CdS to TiO2, mediated by Au-Ag. A strong field is essential for the charge separation at the purposely built interfaces. The band structures of Au₃Ag₇-TiO₂ and Au₃Ag₇-CdS (Fig. 11 (c) and (f)) show that the band structures of both materials certain metallic properties, which benefit the migration of electrons

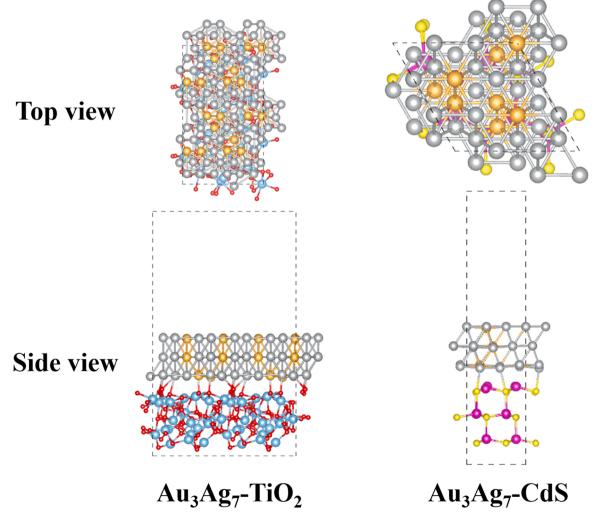


Fig. 10. The optimized structure models of Au₃Ag₇-TiO₂ and Au₃Ag₇-CdS materials.

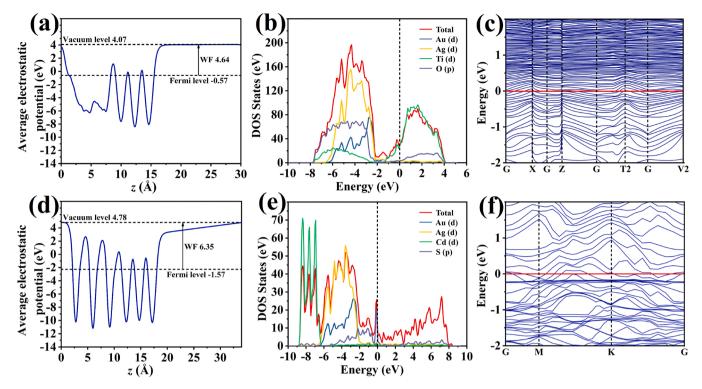


Fig. 11. (a) Calculated Fermi levels and work functions of Au_3Ag_7 -TiO₂. (b) Density of states and partial density of states for Au_3Ag_7 -TiO₂. (c) Band structure for Au_3Ag_7 -TiO₂. (d) Calculated Fermi levels and work functions of Au_3Ag_7 -CdS. (e) Density of states and partial density of states for Au_3Ag_7 -CdS. (f) Band structure for Au_3Ag_7 -CdS.

3.8. Photocatalytic mechanism discussion

The intrinsic mechanism of the CdS/Au-Ag/B-TiO $_2$ photocatalytic system has been systematically analyzed with using FDTD simulations and DFT calculations [4]. The obtained potentials could be converted into normal hydrogen electrode (NHE) potentials by the following formula : $E_{\rm NHE} = E_{\rm AgCl} + E_{\rm 0AgCl} + 0.059$ pH [48]. Hence, the calculated VB edge positions of the CdS, anatase, and TiO $_2$ (B) were estimated to be 1.41, 2.5, and 2.34 V $_{\rm NHE}$, while the corresponding CB positions were calculated to be - 0.99, - 0.7, and - 1.26 V $_{\rm NHE}$ (Fig. 12), which are more negative than the redox potential of CO $_2$ reduction to CO (CO $_2$ /CO, - 0.53 V $_{\rm NHE}$) and CH $_4$ (CO $_2$ /CH $_4$, - 0.24 V $_{\rm NHE}$), implying the photoexcited electrons are negative enough to facilitate the photocatalytic reduction reaction over the CdS/Au-Ag/B-TiO $_2$ [67].

The Au-Ag alloy nanoparticles play the role of electron mediators in this system, promoting the complexation of low-energy electron and low-energy hole phases in the system, thus improving the utilization of high-energy electrons. Meanwhile, the powerful SPR effect of Au-Ag

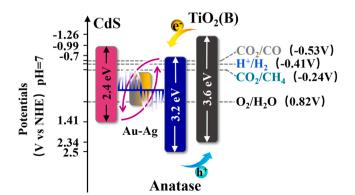


Fig. 12. The intrinsic reaction mechanism of the CdS/Au-Ag/B-TiO $_{\!\!2}$ catalytic system.

alloy promotes this process. Through the RET mechanism, the photoabsorption is not limited by the semiconductor bandgap energy in the full spectrum, and carriers are generated and separated at energy levels above and below the bandgap. Meanwhile, a built-in electric field at the interface between Au₃Ag₇-CdS and Au₃Ag₇-TiO₂ drives the transfer of low-energy electrons from TiO2 to Au-Ag NPs and the low-energy photogenerated holes from CdS to Ag-Au alloy, which are recombined at Au-Ag NPs. Finally, the high-energy photoexcited electrons will stay on the CM of CdS, while the high-energy holes will stay on the VB of TiO₂. The photoexcited charges are well separated, verified by the DFT calculations. The strong electromagnetic field on the Au-Ag NPs will accelerate the transfer of the low-energy electrons and holes. The FDTD simulation confirmed the LEMF mechanism. The CdS/Au-Ag/B-TiO₂ composite structure provides two electric fields to drive the charge separation. One is the local electric field generated by LSPR, and the other is the built-in electric field through the interface. The superposition of these two electric fields can significantly improve the photocatalytic performance of the CdS/Au-Ag/B-TiO2. In addition, given the complexity of the mechanism, the system mechanism in different wavelength cases should be considered. The Au-Ag alloy absorbs visible light and generates an electromagnetic field by LSPR, so the LEMF, RET mechanism, and Z-type heterojunction electron transfer mechanism operate well in the full spectrum and under visible light ($\lambda > 420$ nm) irradiation only. However, the role of Au-Ag alloy in the CdS/Au-Ag/B-TiO2 system is greatly limited under UV irradiation. Under such conditions, the Au-Ag will serve as an electron transfer mediator.

4. Conclusions

In summary, the novel CdS/Au-Ag/B-TiO $_2$ composite heterojunction was successfully prepared with RET mechanism to improve the visible light photocatalytic performance. The established composite heterojunctions (II-type and Z-type) are responsible for the improved charge separation. Meanwhile, the combined characteristics of the photocatalytic Z-type heterojunction system with a solid electronic medium,

bimetallic nanoparticles are used to replace the traditional electronic medium to construct a dual electric field structure with the possible coupling between the electric field at the interface and the plasmon electromagnetic field. Here, two photocatalytic performance tests were performed over the CdS/Au-Ag/B-TiO2: photocatalytic HER reaction and photocatalytic CO₂RR reaction. The sample Cd3Ti2 has the highest photocatalytic property, and the H₂ generation efficiency under the full solar spectrum is 15.97 mmol·h $^{-1}$ ·g $^{-1}$. Its good stability is demonstrated by maintaining 81.3% of performance after 24 h cycles of reactions. The CO_2RR was performed with the CH_4 yield of 14.2 μ mol·h⁻¹·g⁻¹ and the CO yield of 113.9 μ mol·h⁻¹·g⁻¹. High-performance benefits from coordinating the action mechanism between the electric field at the interface and the plasma electromagnetic field coupling. It helps to reduce the random migration of low-energy photo-generated electrons and holes, which improves the directional charge transfer kinetics. The FDTD simulations and DFT calculations help to establish the photocatalytic enhancement mechanism. This work is of great significance for understanding the structure-property-catalysis interplays of plasmonic metalbased heterostructures and realizing the rational design of plasmonenhanced catalysis systems.

CRediT authorship contribution statement

Shengchao Yang and Zhiyong Liu proposed and conceptualized this project. Qi Li and Ruihan Liu designed the experiments. Qi Li carried out the experiments and characterizations. Yufan Huang, Yuwei Liang, and Chunling Hu participated in the experiments or characterizations. Qi Li and Yufan Huang analyzed the results and drew the figures. Qi Li wrote the original draft. Zhiyong Liu, Jichang Liu, Yanlong Tai, and Yongsheng Li reviewed and edited the manuscript. Zhiyong Liu and Shengchao Yang funded this research. Zhiyong Liu, Shengchao Yang, and Yongsheng Li provided the necessary conditions for this research and participated in discussions during data analysis.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

No data was used for the research described in the article.

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Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.apcatb.2023.122700.

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